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# Glass forming ability of ZBLAN glass

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A ZBLAN glass of the following composition: 53 mol.% ZrF4, 20 mol.% BaF2, 4 mol.% LaF3,3 mol.% AIF3, 20 mol.% NaF is studied for predicting its glass forming ability (GFA). The GFA of a ZBLAN glass is evaluated by estimating the Gibbs free energy difference ( $\Delta G$ ) between the undercooled liquid and the corresponding equilibrium solid phases. Other GFA criteria ( $T_{rg}=T_g/T_1$ ,  $T_x/T_1$ ,  $\gamma=T_x/T_g+T_1$ ,  $\Delta T_{xg}=T_x-T_g$ ,  $T_g/T_m$ ) are also calculated. Here, the approximation for  $\Delta G$  considering constant specific heat difference ( $\Delta C_P$ ) over a wide temperature range is calculated for a ZBLAN glass. The so-calculated  $\Delta G$  is compared with the result of other theoretical approaches and experimental data. We can see from the result that the expression with the consideration of constant  $\Delta C_P$  works well in the entire undercooling for ZBLAN glass.

Keywords: ZBLAN glass, Glass forming ability, Gibbs free energy difference ( $\Delta G$ )

# **1** Introduction

A ZBLAN glass: 53 mol.% ZrF4, 20 mol.% BaF2, 4 mol.% LaF3,3 mol.% AIF3, 20 mol.% NaF is a heavy metal-based fluoride system. ZBLAN glass is the most stable fluoride glass known and is most commonly used to make optical fibre (wiki). Superior infrared transmittance is the advantage of ZBLAN over other glasses available, such as silica. It has a broad optical transmission window extending from 0.3 micrometers in the UV to 7 micrometers in the infrared<sup>1</sup>. Fragility and sensitivity to acids are their drawbacks. Formally, fragility reflects the degree to which the temperature dependence of the viscosity (or relaxation time) deviates from Arrhenius behaviour<sup>2</sup>. So it is interesting to study the glass forming ability (GFA) of ZBLAN glass with thermodynamic point of view, which may help to understand the physical origin of this glass. Since complete investigations and detailed descriptions of heavy metal fluoride (HMF) glasses of their properties can be found in the literature<sup>3-12</sup>. The thermodynamic properties of this system are studied earlier<sup>13</sup> but there are not many studies about the thermodynamic properties of fluoride based glass.

It is found that the glass transition temperature Tg is well separated from the crystallization temperature  $T_{x,}$  in fact it is opposite in ZBLAN glass. In the past,

many criteria have been proposed for predicting the glass forming ability (GFA) for metallic glass formers<sup>14-25</sup>.

Here we evaluate the GFA by estimating the Gibbs free energy difference ( $\Delta G$ ) between the undercooled liquid and the corresponding equilibrium solid phases. Other GFA criteria ( $T_{rg}=T_g/T_l$ ,  $T_x/T_l$ ,  $\gamma=T_x/T_g+T_l$ ,  $\Delta T_{xg}=T_x-T_g$ ,  $T_g/T_m$ ) are assessed. In present paper, we have calculated the value of  $\Delta G$  by the expression taking the specific heat difference  $\Delta C_p$  constant without any approximation. It is compared with other expressions and experimental data.

# 2 Theory

The GFA is considered in three different ways namely kinetic, structural and thermodynamic point of view. From the kinetic point of view, viscosity is closely related to the reduced glass transition temperature  $(T_{rg}=T_g/T_l)$  where  $T_g$  is glass transition temperature and  $T_1$  is liquidus temperature<sup>14</sup>. Higher  $T_{rg}$  leads to the higher  $GFA^{26}$ . The  $T_x/T_1$  ratio is also an indicator of the GFA where T<sub>x</sub> is onset crystallization temperature and T<sub>1</sub> is liquidus temperature. The large  $T_x/T_1$  ratio indicates a lower  $R_c$ and therefore the higher GFA<sup>27</sup>. R<sub>c</sub> is a critical cooling rate, which is the minimum cooling rate necessary for keeping the melt amorphous without crystallization. Actually  $R_c$  is a key parameter to evaluate the GFA of a melt but it is difficult to measure precisely. So we have used simple and reliable criteria for the GFA.

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From the structural point of view, higher  $T_g/T_m$  ratio and higher  $\Delta T_{xg}=T_x-T_g^{-17}$  is required for the higher GFA. These are the parameters to gauge the glass stability.

Thermodynamically,  $\Delta G$  is a key parameter for predicting the GFA. Low value of  $\Delta G$  means the driving force for crystallization is low and glass formation becomes easy. In the relation  $\Delta G = \Delta H_f - T\Delta S_f$  for Gibbs free energy,  $\Delta H_f$  and  $\Delta S_f$  are enthalpy of fusion and entropy of fusion. In multi component systems, the number of microscopic states is large leading to the increase in  $\Delta S_f$ , leads to a decrease in  $\Delta G$ . Small value of  $\Delta G$  is favourable for the higher GFA.

# **3** Formulation of Theoretical Expression for Gibbs Free Energy Difference ( $\Delta G$ )

The difference in Gibbs free energy between the liquid and crystalline phases is given by:

$$\Delta G = \Delta H - T \Delta S \qquad \dots (1)$$

where

$$\Delta H = \Delta H_m - \int_T^{T_m} \Delta C_p dT \qquad \dots (2)$$

and

$$\Delta S = \Delta S_m - \int_T^{T_m} \Delta C_p \frac{dT}{T} \qquad \dots (3)$$

where,  $T_m$  is the melting temperature,  $\Delta S_m$  is the entropy of fusion and  $\Delta H_m$  is the enthalpy of fusion.  $\Delta C_p = C_p^l - C_p^x$ , is the difference in specific heats of the liquid and corresponding crystalline phases.  $\Delta S_m$  and  $\Delta H_m$  are related to each other by the following relation:

$$\Delta S_{m} = \Delta H_{m}/T_{m} \qquad \dots (4)$$

Thus, if we have experimental specific heat data for the undercooled liquid and the corresponding equilibrium solids, experimental  $\Delta G$  can be calculated using Eqs (1)-(3). It is difficult to obtain accurately the heat capacity data experimentally because the nature of the undercooled liquid is metastable. So  $\Delta G$ has to be estimated theoretically. The most common expression of the temperature dependence of a linear relationship of  $\Delta C_p$  is given by:

$$\Delta C_{\rm p} = AT + B \qquad \dots (5)$$

where A and B = coefficient for linear variation, T= temperature. Substituting Eq. (5) into Eq. (2) and (3), Eq. (1) can be simplified to:

$$\Delta G = \frac{\Delta H_{\rm m} \Delta T}{T_{\rm m}} - \frac{1}{2} A (\Delta T)^2 + B \left( T \ln \frac{T_{\rm m}}{T} - \Delta T \right) \qquad \dots (6)$$

where  $\Delta T=T_m$ -T To simplify Eq. (6), Thompson & Spaepen<sup>28</sup> used the following approximation:

$$\ln\left(\frac{T_m}{T}\right) = \ln\left(1 + \frac{\Delta T}{\left(T_m + T\right)/2}\right) \cong \frac{2\Delta T}{T_m + T} \qquad \dots (7)$$

and derived the expression:

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$$\Delta G = \frac{\Delta H_m \Delta T}{T_m} \left( \frac{2T}{T_m + T} \right) \qquad \dots (8)$$

The constant  $\Delta C_p$  accounted for the enthalpies of crystallization  $\Delta H_x$  of the metallic glass.

$$\Delta H_{x} = \Delta H_{m} - \Delta C_{p}(T_{m} - T_{x}) \qquad \dots (9)$$

$$\Delta C \mathbf{p} = \gamma \, \frac{\Delta H \, m}{T_{m}} \qquad \dots (10)$$

$$\gamma = \frac{1 - \frac{\Delta H_x}{\Delta H_m}}{1 - \frac{T_x}{T_m}} \qquad \dots (11)$$

Considering this  $\Delta C_p$  value of Eqn. (10) Battezzati *et al.*<sup>29</sup> gives the following expression:

$$\Delta G = \frac{\Delta H_m \Delta T}{T_m} - \gamma \frac{\Delta H_m}{T_m} \left[ \Delta T - T \ln \frac{T_m}{T} \right] \qquad \dots (12)$$

where the value of  $\gamma$  can be calculated by Eq. (11).

The approximation given in Eq. (7) is strictly valid only for small undercooled region,  $\Delta T$ . But, the multi component glass forming systems exhibit a large undercooled regime. Therefore, to account for the wide undercooled region of the multi component metallic alloys, Lad *et al.*<sup>30</sup> obtained the following expression.

$$\Delta G = \frac{\Delta H_m \Delta T}{T_m} \left( 1 - \frac{\Delta T}{2T} \right) \qquad \dots (13)$$

Different approximations for the logarithmic term  $(\tau)$ 

 $\ln\left(\frac{T_m}{T}\right)$  given in Eq. (6) are used by Eq. (8) and

Eq. (13) and it has been found that neither Eq. (8) nor Eq. (13) works for many multi component amorphous alloys.

Considering the Taylor series expansion of Eq. (7) up to second order gives the following approximation:

$$\ln\left(\frac{T_m}{T}\right) \approx \frac{4T\Delta T}{\left(T_m + T\right)^2} \qquad \dots (14)$$

Substituting Eq. (14) in Eq. (6) Lad *et al.*<sup>31</sup> obtained the following equation for the Gibbs free energy difference:

$$\Delta G = \frac{\Delta H_m \Delta T}{T_m} \left( \frac{4T^2}{\left(T + T_m\right)^2} \right) \qquad \dots (15)$$

Most of the bulk glass forming systems shows good result with Eq. (15). But there is some deviation for many systems. It is observed that the specific heat of the undercooled liquid does not vary much with temperature, hence  $\Delta C_p$  is nearly constant. For constant  $\Delta C_p$ , assuming A=0, B=  $\Delta C_p$  =constant then Eq. (6) reduced to:

$$\Delta G = \frac{\Delta H_m \Delta T}{T_m} + \Delta C_p \left( T \ln \frac{T_m}{T} - \Delta T \right) \qquad \dots (16)$$

Following the argument of vanishing  $\Delta S$  at the Kauzmann Temperature  $T_k$ ,

Substituting Eq. (4) in Eq. (3) and  $\Delta S=0$  at  $T=T_k$ , we can get:

$$\Delta C_p = \alpha \frac{\Delta H_m}{T_m} \dots (17)$$

where  $\alpha = 1/\ln(T_m/T_k)$ . Using the Eq. (17) for  $\Delta C_p$  in Eq. (16) Heena *et al.*<sup>32</sup> obtained following expression:

$$\Delta G = \frac{\Delta H_{\rm m}}{T_m} \left[ \alpha T \ln \frac{T_m}{T} + \Delta T \left( 1 - \alpha \right) \right] \qquad \dots (18)$$

They have not taken any approximation to get Eq. (18) except for the assumption related to constant  $\Delta C_{p.}$ 

Thus our expression for  $\Delta G$  (Eq. (18)) considering constant  $\Delta C_p$  fits very well with experimental data for most of the bulk amorphous alloys. Taking the constant value of  $\Delta C_p = \Delta C_p^m$  at the melting point in Eq. (17), have been derived:

$$\alpha = \frac{\Delta C_p^{\ m} T_m}{\Delta H_m} \qquad \dots (19)$$

In this paper,  $\alpha$  has been derived using the Eq. (19) obtained from Eq. (17) taking the constant value of  $\Delta C_p = \Delta C_p^{m}$  at melting point.

However, for such systems, in which the specific heat increases considerably with undercooling,  $\Delta C_p$  at any temperature  $T_k < T < T_m$  in the undercooled region can be expressed as:

$$\Delta C_p = \frac{\Delta C_p^m T_m}{T} \qquad \dots (20)$$

 $\Delta C_p^{m}$  is being the difference in the specific heats at the melting point. Substituting  $\Delta C_p$  from Eq. (20) in Eq. (2) & (3) and using Eq. (1), the  $\Delta G$  values for such systems (Heena *et al.*<sup>32</sup>) is represented by the following equation:

$$\Delta G = \frac{\Delta H_{m} \Delta T}{T_{m}} - \Delta C_{p}^{m} T_{m} \left[ \ln \frac{T_{m}}{T} - \frac{\Delta T}{T_{m}} \right] \qquad \dots (21)$$

#### **4 Result and Discussion**

Figure 1 shows the results of  $\Delta G$  in the under cooled region for ZBLAN glass obtained using the Eq. (18) and Eq.(21) proposed by Heena *et al.*<sup>32</sup>, Battezzati Eq.(12)<sup>29</sup>, Lad-1 Eq.(13)<sup>30</sup>, Lad-2 Eq. (15)<sup>31</sup>. The results for the ZBLAN glass have been compared with the experimental results of Battezzati *et al.*<sup>13</sup>. The parameters used to calculate  $\Delta G$  for this system are given in Table 1.

The excellent agreement of the result obtained using Heena *et al.* given by Eq. (18) with experimental data can be observed and the close matching of the two curves up to large degree of undercooling can be appreciated. Equation (18) does



energy difference $\Delta G^{13}$ .											
System	$T_g(K)$	$T_x(K)$	$T_m(K)$	$T_{l}(K)$	$\Delta H_m$	$C_p^{m}$					
					J/mol	J/mol K					
ZBLAN	545	620	710	820	13961.36	50.23					

not involve any approximation to evaluate the logarithmic term. Also it is derived considering  $\Delta C_p$ to be constant, because the specific heat of undercooled liquid of most of the bulk glass forming systems does not vary much with temperature. In ZBLAN glass, the result of Eq. (18) shows exact matching with experimental curve means the specific heat of the undercooled liquid of ZBLAN glass forming system remains nearly constant with temperature. Heena et al include  $\alpha$  in Eq. (18). The value of  $\alpha$  is 2.55 for ZBLAN system. However, significant deviation is observed at high under cooling with hyperbolic variation of  $\Delta C_p$  if Eq. (21) is used for the evaluation of  $\Delta G$ . The deviation of hyperbolic expression calculated by Eq. (21) is due to the consideration of specific heat varying with under cooling which otherwise appears to be constant. The plot using the Lad-1 approach shows a larger deviation from the experimental curve. The expression of Lad et al., Eq. (13) is obtained using approximation of logarithmic term and done Taylor series expansion up to second order. But it does not give good result for ZBLAN glass. Lad et al. have also proposed another expression given by Eq. (15) in which they have used Thompson-Spaepen approximation of logarithmic term and have done Taylor series expansion up to second order. The Lad-2 approximation works well only at low under cooling. The result using Battezzati et al., Eq. (12) lies below the experimental curve. At a small

undercooling the deviation is small however it is large at large under cooling. Battezzati et al. have taken the value of  $\gamma$  to be ~3.5. The value of  $\gamma$  is calculated 5.9 for ZBLAN system in present paper. It is difficult to evaluate  $\gamma$  correctly because a multi component system crystallizes in multiple steps. It is difficult to choose a particular step which should be taken for the derivation of  $\gamma$ . Besides this,  $T_x$  is heating rate dependent and even  $\Delta H_x$  is not constant for all heating rate. So the value of  $\gamma$  is not unique while  $\alpha$  given by Eq. (19) and used in Eq. (18) for the estimation of  $\Delta G$ is unique. It requires knowledge of  $\Delta C_p^{m}$ ,  $\Delta H_m$  and  $T_m$  Eq. (18) is derived by taking the constant value of  $\Delta C_{\rm p} = \Delta C_{\rm p}^{\rm m}$  at the melting point and the result using Eq.(18) gives very good match with experimental data. Thus, we can say that the value of  $\alpha$  is much appropriate than the value of  $\gamma$  for the ZBLAN glass.

The results of the other GFA parameters are given in Table 2. The results obtained using different methods are compared with the corresponding value for typical bulk metallic glasses (BMGs). The value of reduced glass transition temperature  $T_{rg}=T_g/T_l=0.66$ shows that ZBLAN glass possesses the higher GFA. The large  $T_x/T_1=0.75$  ratio indicates a lower R<sub>c</sub> and therefore a higher GFA. The large super cooled region  $\Delta T_{xg}$ =75 is observed. It indicates that the super cooled liquid has a high resistance to the crystallization. It leads to the higher GFA for ZBLAN glass. Since ZBLAN glass can be cooled more slowly than 1 K/s without noticeable homogeneous nucleation, it has been considered to be as the most stable HMF glass and the most resistant to crystallization under optical fibre pre form-making and drawing condition.<sup>33</sup> However, optical fibre from this material has excellent transmission characteristics in the IR, but the glass is somewhat susceptible to nucleation and crystallization<sup>34</sup>.

We obtain  $\gamma = T_x/T_g + T_l = 0.45$  and  $T_g/_{Tm} = 0.76$ . We obtain the  $\Delta G$  (T<sub>g</sub>) value using Eq. (18) and it is 2.2 kJ/mol while the value of  $\Delta G$  (T<sub>g</sub>) is 1 kJ/mol for the metallic glass former PdCuNiP with the higher GFA.

### 4.1 Calculation of GFA indicators

Table 1 shows the glass transition temperature  $T_g$ , onset crystallization temperature  $T_x$ , offset fusion temperature (liquidus temperature)  $T_l$ , melting temperature  $T_m$ , total heat needed for melting  $\Delta H_m$  and the specific heat difference  $\Delta C_p^{\ m}$  for ZBLAN glass. All of these characteristic temperatures can be easily determined from single DSC measurements.

Table 2 — Different GFA criteria.											
System	$T_{rg} = T_g / T_l$	$T_x/T_1$	$\gamma = T_x/T_g + T_l$	$\Delta T_{xg} = T_x - T_g K$	$T_g/T_m$	$\Delta G(T_g)$ J/mol					
ZBLAN	0.66	0.75	0.45	75	0.76	2200.31					
Typical BMGs [25]	0.503-0.690	0.53-0.78	0.353-0.464	16.3-117	-	-					

Table 2 presents the summary of  $T_{rg}=T_g/T_l$ ,  $T_x/T_l$ ,  $\gamma=T_x/T_g+T_l$ ,  $\Delta T_{xg}=T_x-T_g$ ,  $T_g/T_m$  and  $\Delta G$  based on the data in Table 1.

# **5** Conclusions

The higher value of reduced glass transition temperature  $T_{rg}=T_g/T_l=0.66$  indicates that the melt of ZBLAN glass becomes more viscous and crystallization becomes difficult. Thus the GFA enhances. This system possesses a large  $T_x/T_l=0.75$  could have a lower R<sub>c</sub> therefore a higher GFA. Large supercooled liquid region  $\Delta T_{xg}=T_x$ -  $T_g=75$  K is an indication of the devitrification tendency of ZBLAN glass upon heating above  $T_g$ . The lower value of  $\Delta G$  ( $T_g$ ) =2.2 kJ/mol for ZBLAN glass leads to the higher GFA. The calculation of Gibbs free energy difference  $\Delta G$  for ZBLAN glass is done by different expressions. The expression with the consideration of constant  $\Delta C_p$  without taking any approximation provides fairly accurate result matching with experimental data.

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