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# Systematics of alpha decay hindrance factors in doubly-even nuclei

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In present work, we have calculated the hindrance factors of 182 even-even alpha emitters using Preston's formulation of alpha decay probabilities and presented HFs systematics of 1<sup>-</sup> and 3<sup>-</sup> states in reflection asymmetric even-even quadruple-octupole deformed nuclides (A~216-230). The calculated HFs of both 1<sup>-</sup> and 3<sup>-</sup> states decrease with reduction in neutron number and this decrease is attributed to onset of intrinsic reflection asymmetry. There is a trend reversal for 1<sup>-</sup> states at N=132 ( $^{218}$ Ra) and N=134 ( $^{220}$ Rn), which might be a possible indication of departure from static octupole deformation. Similarly, HFs systematics is discussed for  $^{224-230}$ Th and  $^{232-236}$ U isotopic chains along with 2<sup>+</sup> states observed in daughter nuclei in N=132-146 isotonic chains.

Keywords: Even-even alpha emitters, Hindrance factors, Reflection asymmetry

#### **1** Introduction

Alpha hindrance factor (HF), the ratio of experimental to theoretical partial half-lives of alpha transitions, is found to be an important tool in extracting nuclear structure information<sup>1-4</sup>. Various theoretical techniques have been developed in order to understand the alpha-decay process and hence to calculate the penetrability of alpha particles through a barrier<sup>1,5</sup>. The alpha transitions for which HF lies between 1 and 4, called favored transitions and take place between nuclear states having similar configurations and hence it is promising to ascertain both  $J^{\pi}$  and nucleonic configurations assignments for a given daughter (parent) state if those of parent (daughter) are known<sup>6</sup>. Similarly, the alpha HFs quantifies the correlation between nuclear wave functions of the initial state of parent and final state of daughter nuclei; larger wave function's overlap gives a lower HF<sup>3</sup>. The systematics of alpha-decay HFs can be used to deduce a variety of quantities like total alpha branching ratio, intensities of unobserved alpha groups and excitation energy of level in daughter nucleus<sup>6</sup>. In the present study, the spin-independent part of Preston's equations<sup>1</sup> have been used for the calculations of alpha decay probabilities. This formalism contains radius of the daughter nuclide,  $r_0$ , as a free parameter. By setting HF=1 for the groundstate to ground-state alpha branch for an even-even nuclide<sup>2</sup>, the radius parameter for the daughter nuclide can be deduced<sup>7</sup> that can be used to deduce alpha hindrance factors for alpha-fed excited states in eveneven nuclides. We have calculated HFs of 182 eveneven alpha emitters in the framework of Preston's formulation<sup>1</sup> by using ALPHAD\_RadD program<sup>8</sup> and present the systematics of alpha HFs with daughter neutron number in reflection asymmetric (RA) mass region (A=216-230) i.e. for quadruple-octupole deformed nuclei. Additionally, the systematics of HFs for 2<sup>+</sup> states observed in N=132-146 isotonic chains is also presented

# 2 Methodology

In the present study, a well established Preston's spin-independent formulation<sup>1</sup> with only essential steps described here, has been used. In this formulation, a preformed alpha particle is considered to be moving inside a nucleus having rectangular potential field of depth of  $V_0$ ;  $V_0 = \text{constant}$  for  $r < r_0$ , where r is the distance from the center of the product nucleus and  $r_0$  is the radius of the product nucleus. The field beyond effective nuclear radius was assumed to be generated by a Coulomb potential  $(2Ze^2/r, \text{ where } Z \text{ is the atomic number of product nucleus and e is the elementary charge) between alpha particle and daughter nucleus. The solution of the time dependent Schrodinger equation is assumed to have a form<sup>1</sup>:$ 

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 $u = \psi(x, y, z) \exp(-iEt / \hbar)$ 

The wave-function  $\psi$  should obey following boundary conditions (a) at r = 0 at  $\psi$  is finite, (b) at  $r = r_0$ ,  $\psi$  and  $d\psi/dr$  are continuous, (c)  $\psi$  represents as outgoing wave for  $r > r_0$ . The Schrödinger equation for alpha particle in afore said potential field can be written as:

$$\frac{d^2 X_l}{dr^2} + \frac{2m}{\hbar^2} \left[ E - V - \frac{\hbar^2 l \left( l + 1 \right)}{2mr^2} \right] X_l = 0$$

where,  $E = E_{\alpha} - \frac{1}{2}i\hbar\lambda$ ;  $E_{\alpha}$  is the energy of  $\alpha$ -particle

and  $\lambda$  the time constant; are complex eigen values. In the interior of the nucleus with V=U, we have

$$X_{l}^{i} = \left\{\frac{2m(E-U)}{\hbar^{2}}\right\}^{\frac{1}{4}} r^{1/2} J_{l+\frac{1}{2}} \left[\left\{\frac{2m(E-U)}{\hbar^{2}}\right\}^{\frac{1}{2}} r\right]$$

where, J denotes Bessel functions and superscript i refer to interior of the nucleus. The solution  $X_i^i$  represents a standing wave and imaginary part of E related to the leak of alpha particle through the potential barrier.

#### **3 Input Parameters**

In the present work, the daughter radius parameter  $(r_0)$  is the main input used to calculate the HFs of alpha-fed excited states in even-even nuclides. In order to deduce  $r_0$  a set of experimental quantities

such as  $Q_{\alpha}$  value, parent nuclide's half-life (T<sub>1/2</sub>), total alpha–decay branching ratios (% $\alpha$ ), and alpha intensities (I<sub>a</sub>) are used. The  $Q_{\alpha}$  energies are taken from recent atomic mass evaluation of M. Wang et al.<sup>9</sup>, total alpha-decay branching ratios and half-lives of parent nuclides are taken from the ENSDF database<sup>10</sup> supplemented by recent data from literature.

In order to calculate HFs of various excited states, the recently developed ENSDF analysis code namely, ALPHAD\_RadD<sup>8</sup>, which is based on Preston's equations for alpha decay transition probabilities<sup>1</sup>, has been used. This program can also be used to deduce HFs in odd-A and odd-odd nuclei by employing recently evaluated<sup>7</sup> list of even-even daughter radius parameters.

## **4 Results and Discussion**

In Reflection Asymmetric even-even quadrupleoctupole deformed nuclides (A~216-230), two separate bands with opposite parity i.e.  $I^{\pi}=0^+$ ,  $2^+$ ,  $4^+$ ... and  $I^{\pi}=1^-$ ,  $3^-$ ,  $5^-$ ... are generally observed<sup>4</sup>. In present paper, we studied the systematics of alpha HFs for 1<sup>-</sup> and 3<sup>-</sup> states with daughter neutron number in above said RA mass region (A~216-230). The results of HF systematics of 1<sup>-</sup> and 3<sup>-</sup> states observed in <sup>218-226</sup>Ra isotopic chains with daughter neutron number are presented in Fig.1(a). From this figure, it is clear that, the HFs of both 1<sup>-</sup> and 3<sup>-</sup> states smoothly decreasing with reduction in neutron number. This decrease of HFs is attributed to onset of

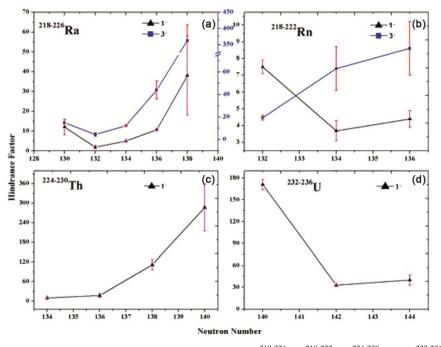


Fig. 1(a-d) — HF's systematics with daughter neutron number for <sup>218-226</sup>Ra, <sup>218-222</sup>Rn, <sup>224-230</sup>Th and <sup>232-236</sup>U nuclides.

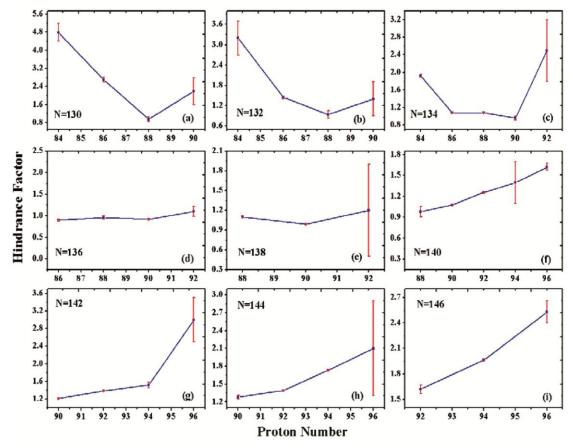


Fig. 2(a-i) — HF's systematics of 2<sup>+</sup> states observed in daughter nuclei in N=132-146 isotonic chains.

static Quadruple-Octupole deformation, but there is a trend reversal at N=132 ( $^{218}$ Ra), which might be a possible indication of departure from static Quadruple-Octupole deformation<sup>3</sup>.

Similar trend reversal in HFs systematics of 1<sup>-</sup> state at N = 134 is also observed for <sup>218-222</sup>Rn nuclides (Fig.1 (b)), but such reversal could not be observed in <sup>224-230</sup>Th (Fig.1(c)) as experimental data for 1<sup>-</sup> state in lower mass region is not accessible. The HF systematics of 1<sup>-</sup> state in <sup>232-236</sup>U nuclides (Fig.1(d)) shows a minimum at N=142 and may be due to similar shape transition as discussed for Ra and Rn isotopic chains.

Additionally, we have also presented the systematics of HFs of  $2^+$  states observed in N=132-146 isotonic chains as shown in Fig. 2(a-i). From Fig. 2(a-c), a smooth decrease in HFs of  $2^+$  states with increase in proton number is observed till Z=88 in N=130, 132 and till Z=90 in N=134 isotopic chain. This smooth decrease in HFs indicates more probable alpha penetration through coulomb barrier and hence could be attributed to decrease in stability beyond Z=82 shell closure. The abrupt increase of HF at

Z=90 in N=130, 132 and at Z=92 in N=134 isotopic chain is still an open problem. Although there is increase in HFs in other isotonic chains (N=136-146) as shown in Fig. 2(d-i), but minima corresponding to certain Z number could not be identified due to inaccessibility of experimental data for these isotonic chains. On the basis of above systematics (Fig. 2(a-i)), we suggest that all the experimentally observed alpha decays branches (0<sup>+</sup> to 2<sup>+</sup>) are favored.

#### **5** Conclusions

The HFs systematics of 1<sup>-</sup> and 3<sup>-</sup> states observed in RA mass region is presented. The HFs of both 1<sup>-</sup> and 3<sup>-</sup> states decrease with reduction in neutron number and this decrease of HFs could be attributed to onset of intrinsic reflection asymmetry, but there is a trend reversal for 1<sup>-</sup> state at N=132 ( $^{218}$ Ra) and N=134 ( $^{220}$ Rn), which might be a possible indication of departure from static octupole deformation. A smooth decrease in HFs of 2<sup>+</sup> states with proton number is observed N=130, 132 and N=134 isotonic chains, which indicates the enhanced alpha penetration probability beyond Z=82 shell closure.

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